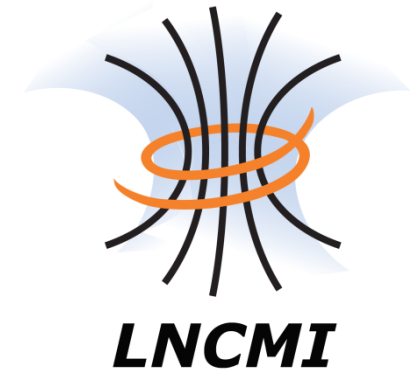


Casimir momentum in crossed electromagnetic fields. QED correction to Abraham force?

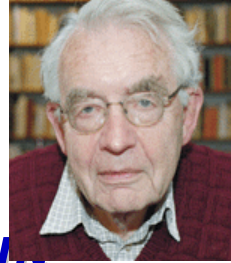


Bart van Tiggelen *and* ***Geert Rikken***

- Sébastien Kawka (Ph.D Grenoble → ENS Pisa)
- James Babington (postdoc ANR Grenoble)

Casimir Workshop ; Leiden , March 2012

Casimir energy

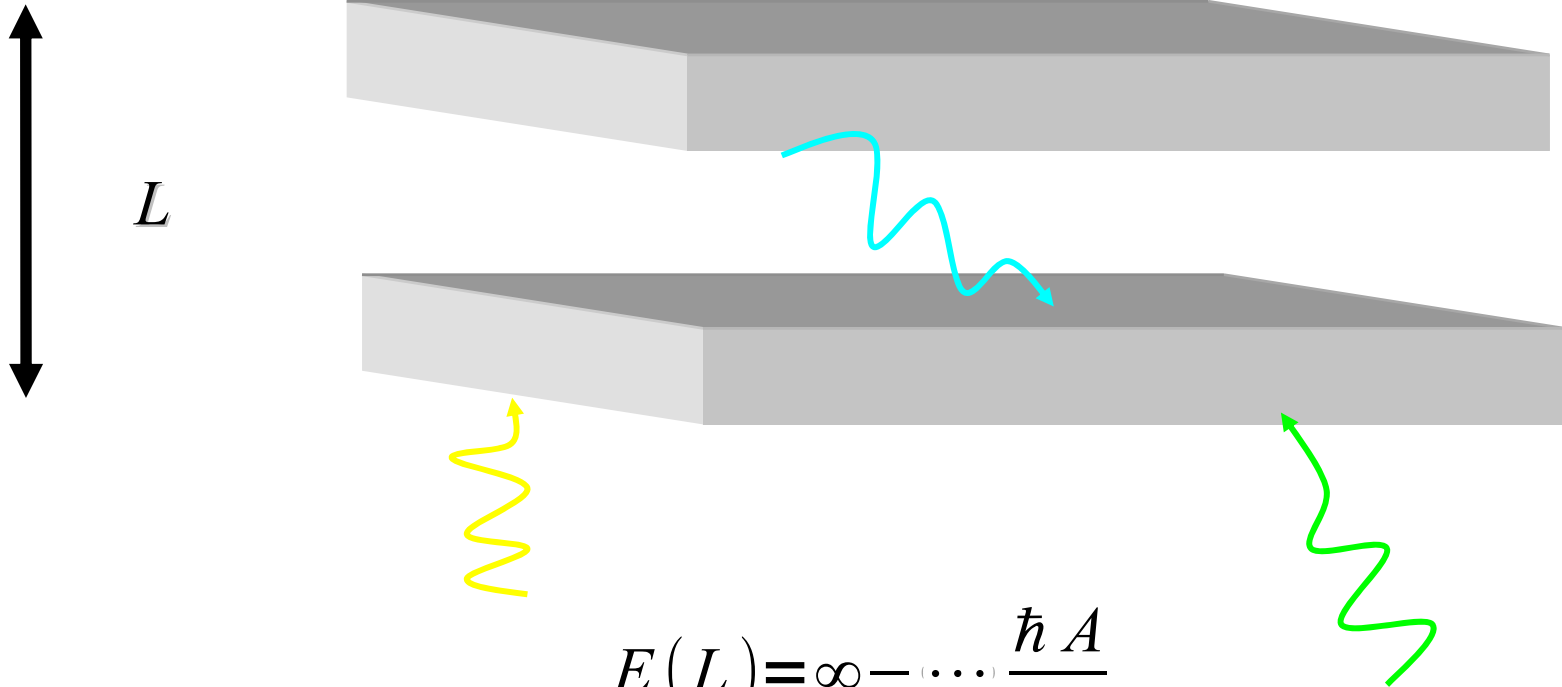


1. **Isotropic radiation with power spectrum ω^3 is Lorentz-invariant (Einstein, 1917);**

$$\frac{1}{c_0^3} \int d\omega \int d\hat{\mathbf{k}} \omega^2 \left[\frac{1}{2} \hbar \omega, \frac{1}{2} \hbar \mathbf{k} \right] = [\rho_{casi}, 0] \quad \text{is an invariant four vector}$$

1. **Van der Waals force $1/r^6$ (London, 1930)**
2. **Relation to Cosmological constant (Pauli, 1934, Davies, 1984)**
3. **Casimir Polder Force $1/r^7$ (1947)**
4. **Attraction between metallic plates (Casimir, 1948)**
5. **Lifshitz theory for dielectric media (Lifshitz, 1956, Dzyalovich, 1961)**
1. **Observation of Casimir effect (Spurnaay, 1958, Lamoureux (5%), 1997), Chan et al, (1%), 2001)**
1. **Stability of the electron (Casimir, 1956, Boyer, 1968)**
2. **Unruh effect & Hawking radiation (Hawking 1974, Unruh 1976)**
3. **Bag model for Hadrons (Jaffe et al, 1974)**
4. **Sign of the Cosmological constant (Weinberg, ... 1983)**
5. **Sonoluminescence (Schwinger, 1993, Eberlein, 1996)**
6. **Quantum friction and sheering the quantum vacuum (Pendry, 1998)**
7. **Casimir momentum in magneto-electric media (Feigel, 2004)**

The Casimir effect....

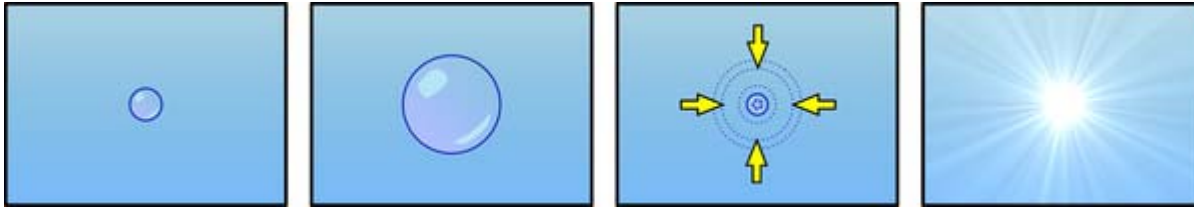


$$E(L) = \infty - \dots - \frac{\hbar A}{L^3}$$

$$F(L) = -\frac{\partial E}{\partial L} = -3(\dots) \frac{\hbar A}{L^4} \hat{L} \quad \text{Negative pressure}$$

No momentum exchange between matter and radiation

UV catastrophe in sonoluminescence (> 1934)



Schwinger (1993)

$$\Delta E(\text{bubble}) = \int d^3 r \left\{ \int d^3 k \frac{1}{2} \hbar \omega_k(\text{bubble in water}) - \int d^3 k \frac{1}{2} \hbar \omega_k(\text{water no bubble}) \right\}$$

$$\approx \frac{\hbar a^3 \omega_c^4}{c_0^3} \left(1 - \frac{1}{\sqrt{\epsilon}} \right) \approx 10 \text{ MeV}$$



cut-off in the UV ?

Dimensional regularisation?



$$\Delta E(\text{bubble}) = \frac{23}{1536\pi} \frac{(\epsilon - 1)^2 c_0}{a}$$

$$\approx 0.001 \text{ eV}$$



Identity of the van der Waals Force and the Casimir Effect and the Irrelevance of These Phenomena to Sonoluminescence

Iver Brevik*

Division of Applied Mechanics, Norwegian University of Science and Technology, N-7491 Trondheim, Norway

Valery N. Marachevsky[†]

Department of Theoretical Physics, St. Petersburg University, 198 904 St. Petersburg, Russia

Kimball A. Milton[‡]

Department of Physics and Astronomy, The University of Oklahoma, Norman, Oklahoma 73019

(Received 12 October 1998)

« Momentum from Nothing »

VOLUME 92, NUMBER 2

PHYSICAL REVIEW LETTERS

week ending
16 JANUARY 2004

Quantum Vacuum Contribution to the Momentum of Dielectric Media

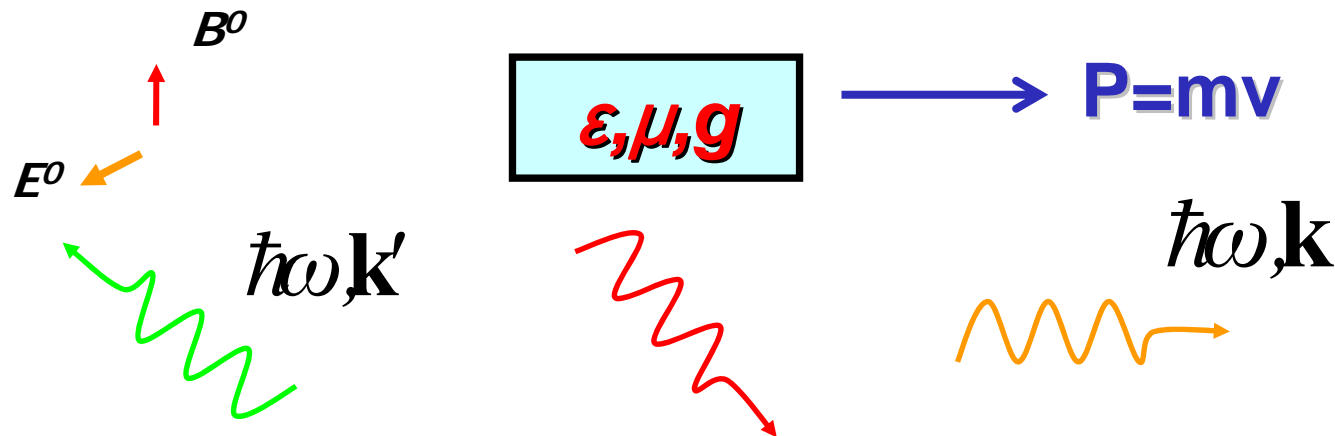
A. Feigel*

Department of Complex Systems, Weizmann Institute of Science, Rehovot 76100, Israel
(Received 3 February 2003; published 16 January 2004)

Momentum transfer between matter and electromagnetic field is analyzed. The related equations of motion and conservation laws are derived using relativistic formalism. Their correspondence to various, at first sight self-contradicting, experimental data (the so-called Abraham-Minkowski controversy) is demonstrated. A new, Casimir-like, quantum phenomenon is predicted: contribution of vacuum fluctuations to the motion of dielectric liquids in crossed electric and magnetic fields. Velocities of about 50 nm/s can be expected due to the contribution of high frequency vacuum modes. The proposed phenomenon could be used in the future as an investigating tool for zero fluctuations. Other possible applications lie in fields of microfluidics or precise positioning of micro-objects, e.g., cold atoms or molecules.

DOI: 10.1103/PhysRevLett.92.020404

PACS numbers: 03.50.De, 42.50.Nn, 42.50.Vk



Bi-anisotropic Media

$$\mathbf{D}(\omega) = \varepsilon(\omega)\mathbf{E}(\omega) + \mathbf{g}(\omega) \cdot \mathbf{B}(\omega)$$

$$\mathbf{H}(\omega) = \mathbf{g}^T(\omega) \cdot \mathbf{E}(\omega) + \mu(\omega)^{-1} \mathbf{B}(\omega)$$

Fresnel dispersion law

$$\det \left(\varepsilon \frac{\omega^2}{c_0^2} - k^2 + \mathbf{k}\mathbf{k} - \frac{\omega}{c_0} \mathbf{g} \cdot (\boldsymbol{\varepsilon} \cdot \mathbf{k}) + \frac{\omega}{c_0} (\boldsymbol{\varepsilon} \cdot \mathbf{k}) \cdot \mathbf{g}^* \right) = 0$$

$$g_{ij}(\omega) = i\omega g \delta_{ij}$$

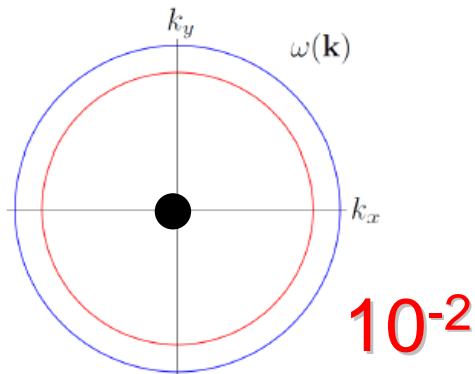
Rotatory power

$$g_{ij}(\omega) = (1 - \varepsilon) \varepsilon_{ijl} \frac{v_l}{c_0}$$

Fizeau effect

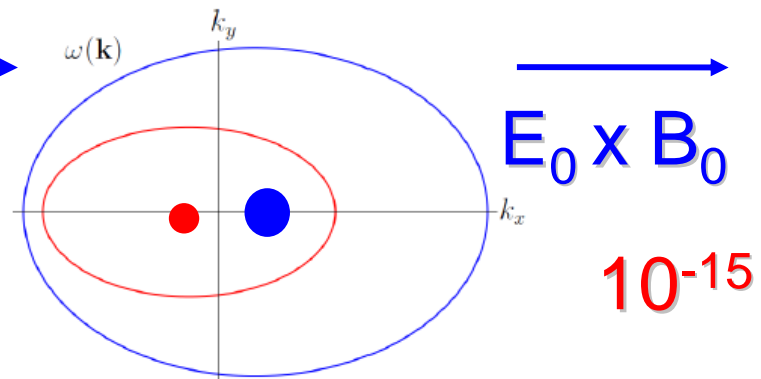
$$g_{ij}(\omega) = g (E_i^0 B_j^0 - B_i^0 E_j^0)$$

Magneto-electric birefringence



\mathbf{v}

10^{-8}




phenomenological continuum theory

$$\partial_t(\rho \mathbf{v} + \varepsilon_0 \mathbf{E} \times \mathbf{B}) = -\nabla \cdot \mathbf{T}^0$$

$$T_{ij}^0 = \frac{1}{8\pi} \left(\varepsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right) \delta_{ij} - \frac{1}{4\pi} \left(\varepsilon_0 E_i E_j + \frac{1}{\mu_0} B_i B_j \right)$$

Observed in X-ray



$$\langle 0 | \frac{\mathbf{E} \times \mathbf{B}}{4\pi c_0} | 0 \rangle \propto \begin{cases} \frac{1}{c_0} \int d^3 \mathbf{k} \frac{1}{2} \hbar \omega_k \times g(\omega) \mathbf{E}_0 \times \mathbf{B}_0 = \frac{2}{3} \frac{\hbar \omega_c^4}{\pi^3 c_0^4} g \mathbf{E}_0 \times \mathbf{B}_0 \\ \frac{1}{c_0} \int d^3 \mathbf{k} \frac{1}{2} \hbar \omega_k \times [\varepsilon(\omega) - 1] \frac{\mathbf{v}}{c_0} = \rho_{casi} \mathbf{v} \end{cases}$$

Photonic momentum in dielectric media?

→ classical « Abraham » contribution already controversial

UV catastrophe of vacuum energy ?

Lorentz invariance of quantum vacuum?

Inertia of quantum vacuum?

The Abraham Force

Macroscopic
Maxwell

$$\rightarrow \partial_t \mathbf{G}_M + \nabla \cdot \mathbf{T} = -\mathbf{f}$$

$$\mathbf{G}_M = \mathbf{D} \times \mathbf{B}$$

$$\mathbf{f} = -E^2 \nabla \varepsilon - H^2 \nabla \mu$$

Minkowski

The Abraham Force

$$\mathbf{G}_A = \varepsilon_0 \mu_0 \mathbf{E} \times \mathbf{H}$$

$$= \frac{1}{c_0} \mathbf{S}$$

$$\partial_t \mathbf{G}_A + \nabla \cdot \mathbf{T} = -\mathbf{f} - \varepsilon_0 (\varepsilon_r - 1 / \mu_r) \partial_t (\mathbf{E} \times \mathbf{B}) \quad \text{Abraham}$$

$$\mathbf{G}_N = \varepsilon_0 \mathbf{E} \times \mathbf{B}$$

$$= \mathbf{G}_0$$

$$\partial_t \mathbf{G}_A + \nabla \cdot \mathbf{T} = -\mathbf{f} - \varepsilon_0 (\varepsilon_r - 1) \partial_t (\mathbf{E} \times \mathbf{B}) \quad \text{Nelson}$$

$$\dots = -\mathbf{f} - \frac{2}{5} [\partial_t \mathbf{E}_0(t)] \times \mathbf{B}_0(t) - \mathbf{E}_0(t) \times [\partial_t \mathbf{B}_0(t)]$$

Peierls

exp

$$\mathbf{F} = (\varepsilon - 1) V \partial_t \mathbf{E}_0(t) \times \mathbf{B}_0(t)$$

Walker & Walker,
Nature 1976

theo

**Abraham momentum = kinetic momentum,
Minkowski momentum = conjugate momentum**

Barnett
(PRL 2010):

Nelson momentum = pseudo momentum

Nelson
(PRA 1991)

The Abraham Force (our version)

**Macroscopic
Maxwell**



$$\partial_t \mathbf{G} + \nabla \cdot \mathbf{T} = -\mathbf{f}$$

$$\mathbf{G} = \mathbf{D} \times \mathbf{B}$$

$$\mathbf{f} = -E^2 \nabla \varepsilon - H^2 \nabla \mu$$

**Maxwell-Lorentz force
on induced polarization
and current**



$$\partial_t \rho \mathbf{v} + \nabla \cdot \mathbf{U} = \mathbf{f} + \partial_t (\mathbf{P} \times \mathbf{B})$$



**Microscopic
Maxwell**



$$\partial_t (\rho \mathbf{v} + \varepsilon_0 \mathbf{E} \times \mathbf{B}) + \nabla \cdot \mathbf{T}_0 = 0$$

The Abraham Force (our version)

Macroscopic
Maxwell



$$\partial_t \mathbf{G} + \nabla \cdot \mathbf{T} = -\mathbf{f} - \partial_t (\mathbf{P} \times \mathbf{B})$$
$$\mathbf{G} = \mathbf{D} \times \mathbf{B}$$

$$\mathbf{f} = -E^2 \nabla \varepsilon - H^2 \nabla \mu$$

Maxwell-Lorentz force
on induced polarization
and current



$$\partial_t \rho \mathbf{v} + \nabla \cdot \mathbf{U} = \mathbf{f} + \partial_t (\mathbf{P} \times \mathbf{B})$$

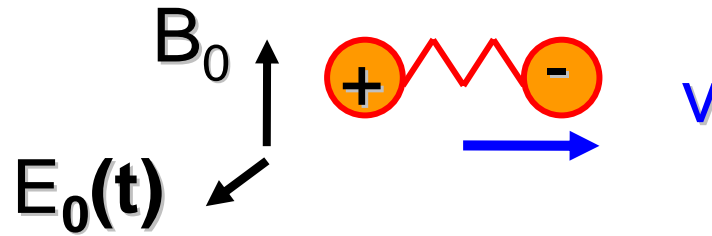


Microscopic
Maxwell



$$\partial_t (\rho \mathbf{v} + \varepsilon_0 \mathbf{E} \times \mathbf{B}) + \nabla \cdot \mathbf{T}_0 = 0$$

Classical Abraham momentum in crossed EM fields



$$\mathbf{r}_{1,2} = \mathbf{R} \pm \frac{1}{2}(\mathbf{x})$$

$$m\ddot{\mathbf{r}}_1 = +q\mathbf{E}(t) + q\dot{\mathbf{r}}_1 \times \mathbf{B} + \mathbf{f}(r_{12})$$

$$2m\dot{\mathbf{R}} + q\mathbf{x} \times \mathbf{B} = \text{constant} = 0$$

$$m\ddot{\mathbf{r}}_2 = -q\mathbf{E}(t) - q\dot{\mathbf{r}}_2 \times \mathbf{B} - \mathbf{f}(r_{12})$$

$$m\ddot{\mathbf{x}} = 2q\mathbf{E}(t) + 2q\dot{\mathbf{R}} \times \mathbf{B} - m\omega_0^2 \mathbf{x} \approx 0$$

$$2m\dot{\mathbf{R}} = \frac{q^2 / m}{\omega_0^2} \mathbf{E}_0(t) \times \mathbf{B}_0$$

No controversy exists in microscopic description
Consistent with Abrahams and Nelson version

The UV catastrophe is real in macroscopic 'description

Free electron (electric dipole)

$$\varepsilon(\omega) - 1 = -\frac{\omega_P^2}{\omega^2} \quad \rho_{casi} = \frac{\hbar}{c_0^3} \int_0^\infty d\omega \omega^3 \frac{\omega_p^2}{\omega^2} = \infty$$

$$g_{ME}(\omega) = -\frac{e^4 m_\Delta^2}{\omega_0^2 m^3 M^2} \left[-\frac{\omega^2 + \omega_0^2}{(\omega_0^2 - \omega^2)^2} (\mathbf{E}_i^0 \mathbf{B}_j^0 - (\mathbf{E}^0 \cdot \mathbf{B}^0) \delta_{ij}) \right. \\ \left. + \frac{1}{\omega_0^2 - \omega^2} (\mathbf{E}_i^0 \mathbf{B}_j^0 - \frac{1}{4} \mathbf{E}_j^0 \mathbf{B}_i^0 - \frac{1}{4} (\mathbf{E}^0 \cdot \mathbf{B}^0) \delta_{ij}) \right].$$

magnetic dipole

Electric quadrupole

$$P_{casi} = \frac{\hbar}{c_0^3} \int dr \int_0^\infty d\omega \omega^3 g(\omega) E_0 \times B_0 = \infty$$

Rizzo et al, 2003-2009, Babington & BAvT, EPJD 2011

$$P_{casi} = \frac{\hbar c_0 g(\varepsilon(0) - 1)}{a} E_0 \times B_0 ?$$

Dimensional regularization for object of size a?

BAvT EPJD 2009

Casimir momentum, if infinite, is Lorentz invariant

$$L(\mathbf{E}, \mathbf{B}) = \frac{1}{2} (\mathbf{E}^2 - \mathbf{B}^2) + 2\nu (\mathbf{E}^2 - \mathbf{B}^2)^2 + \frac{\nu}{2} (\mathbf{E} \cdot \mathbf{B})^2$$

$$\begin{aligned} \mathbf{E} &= \mathbf{E}_0 + \mathbf{E}(\omega) \\ \mathbf{B} &= \mathbf{B}_0 + \mathbf{B}(\omega) \end{aligned}$$

← *Bi-anisotropic
Lorentz-invariant vacuum*

*Fluctuation-
Dissipation*

$$\langle 0 | E_i(\mathbf{r}, \omega) E_j^*(\mathbf{r}', \omega') | 0 \rangle = -2\hbar\omega^2 \text{Im}G_{ij}(\mathbf{r}, \mathbf{r}', \omega) \times 2\pi\delta(\omega - \omega')$$



$$\langle 0 | \frac{c_0}{4\pi} \mathbf{E}^* \times \mathbf{H} | 0 \rangle = 0$$

Zero energy flow

$$\langle 0 | \frac{\mathbf{E}^* \times \mathbf{B}}{4\pi c_0} | 0 \rangle = -\frac{4}{3} \nu K \mathbf{E}_0 \times \mathbf{B}_0$$

infinite momentum density

$$K = \frac{1}{(2\pi)^3} \frac{1}{2} \hbar \int_0^\infty d\omega \int_{4\pi} d\Omega \rho_0(\omega, \Omega)$$

Lorentz scalar

Ex: Helium

$$E_0=450 \text{ V/mm}; B_0=1 \text{ T}$$

$$\alpha(0)=0.22 \cdot 10^{-40} \text{ Cm}^2/\text{V} \quad (16.6a_0^3)$$

$$\rho=0.17 \text{ kg/m}^3 \text{ (room } T)$$

$$g=0.017 \cdot 10^{-22} \text{ m/VT}$$

(SI units)

$$v_{\text{abr}} = \frac{\epsilon_0 \alpha(0) EB}{2m_p} \approx 0.3 \text{ nm/sec}$$

Classical abraham
force

$$F_{\text{abr}} \approx 7 \cdot 10^{-32} \text{ N}$$

$$v_{\text{Feigel}} = \frac{\pi}{4} \frac{h}{\rho \lambda_c^4} gEB \approx 0.02 \text{ nm/sec}$$

$$N_{\text{at}} F_{\text{abr}} \propto 10^{-13} \text{ N}$$

Semi-classical QED with cut-off
0.1 nm (Feigel)

$$v_{\text{QED}} \propto v_{\text{abr}} \times (\alpha)^2 \approx 0.0002 \text{ nm/sec}$$

Rigorous QED (Kawka, 2010,2012)

Ex: Helium

$$E_0=450 \text{ V/mm}; B_0=1 \text{ T}$$

$$\alpha(0)=0.22 \cdot 10^{-40} \text{ Cm}^2/\text{V} \quad (16.6a_0^3)$$

$$\rho=0.17 \text{ kg/m}^3 \text{ (room } T)$$

$$g=0.017 \cdot 10^{-22} \text{ m/VT}$$

(SI units)

$$v_{\text{abr}} = \frac{\epsilon_0 \alpha(0) EB}{2m_p} \approx 0.3 \text{ nm/sec}$$

Classical Abraham

Force

$$F_{\text{abr}} \approx 7 \cdot 10^{-32} \text{ N}$$

$$v_{\text{Feigel}} = \frac{\pi}{4} \frac{h}{\rho \lambda_c^4} gEB \approx 0.02 \text{ nm/sec}$$

$$N_{\text{at}} F_{\text{abr}} \propto 10^{-13} \text{ N}$$

Semi-classical QED with cut-off
0.1 nm (Feigel)

$$v_{\text{QED}} \propto v_{\text{abr}} \times (\mathbf{Z?}\alpha)^2 \approx 0.001 \text{ nm/sec}$$

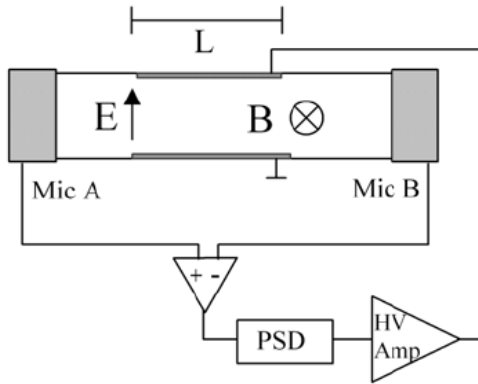
Rigorous QED (Kawka, 2010,2012)

$$\frac{dp}{dt} = \alpha(0) \frac{dE}{dt} \times B$$

Abraham force

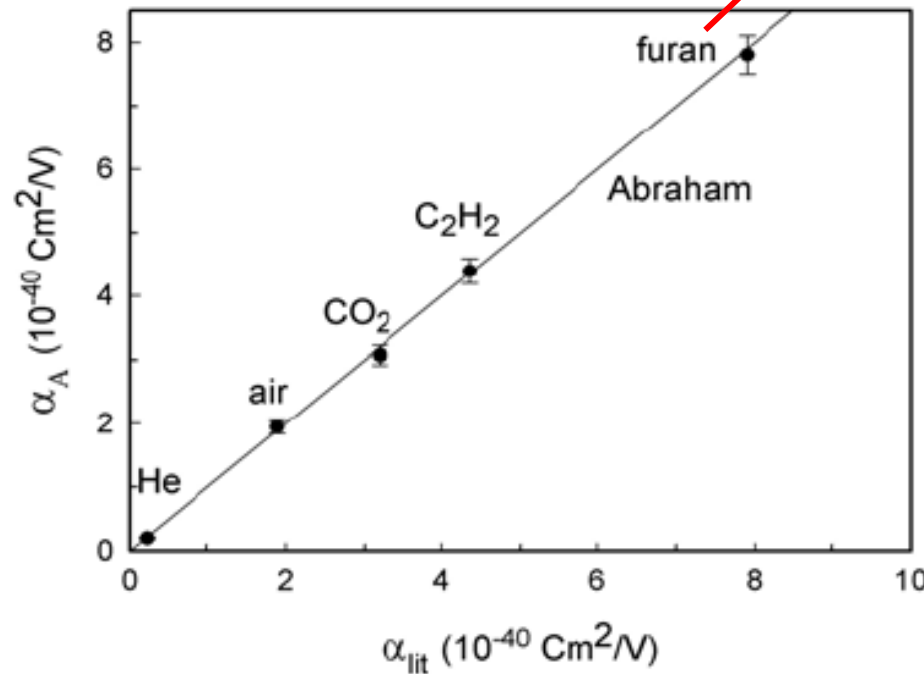
Acoustic pressure

$$P(\omega) = P_0 + \alpha(0) \times E \times B \times \omega \times \cos \omega t \times n \times L$$



V = 8 nm/sec ± 0.8
Feigel correction: 2 nm/sec

P/(EB)



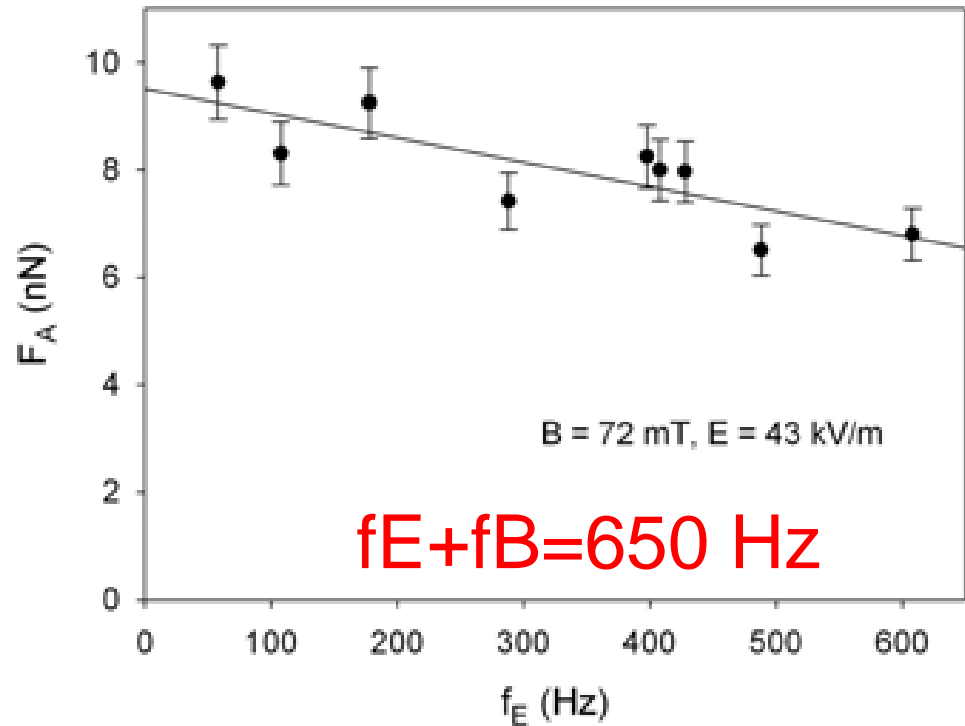
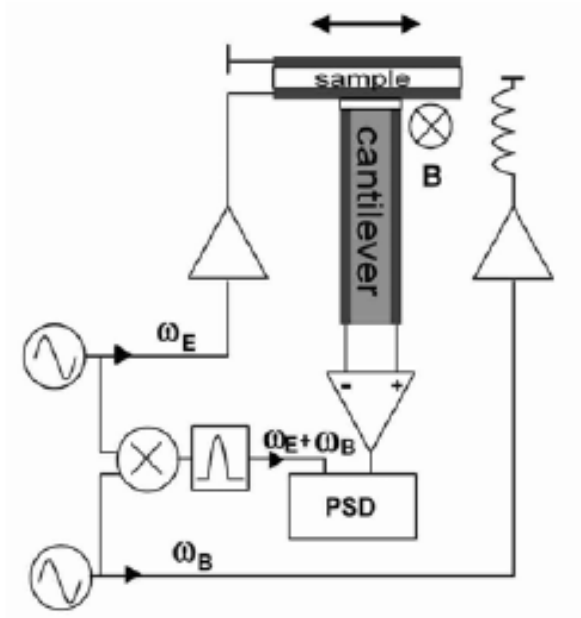
E=450 V/mm;
B=1 T;
f= 7.6 kHz

$\alpha(0)$



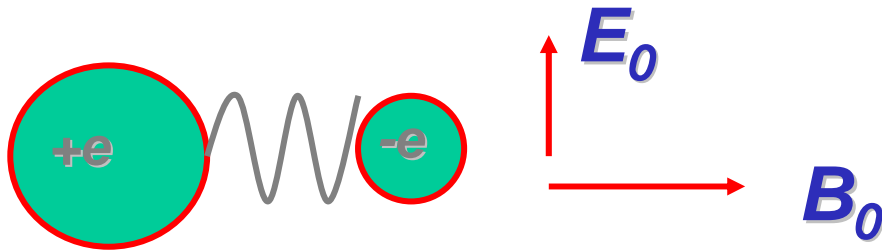
$$\frac{dp}{dt} = \epsilon_0 (\epsilon_r - 1) \frac{d}{dt} (E \times B)$$

$$\epsilon_r = 1.7 \cdot 10^5 \quad Y5V \text{ ceramic}$$



Casimir momentum: 1/6

QED of atom in crossed fields



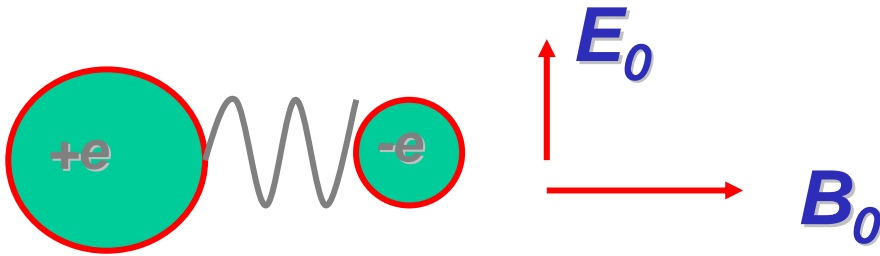
$$\mathbf{A}_0 = \frac{1}{2} \mathbf{B}_0 \times \mathbf{r} \quad \phi = -\mathbf{E}_0 \cdot \mathbf{r}$$

Coulomb Gauge

$$H = \frac{1}{2m_1} (\mathbf{p}_1 - e\mathbf{A}_0(\mathbf{r}_1) - e\mathbf{A}(\mathbf{r}_1))^2 + \frac{1}{2m_2} (\mathbf{p}_2 + e\mathbf{A}_0(\mathbf{r}_2) + e\mathbf{A}(\mathbf{r}_2))^2$$
$$+ e\mathbf{E}_0 \cdot \mathbf{r}_{21} + V(r_{12})$$
$$+ \sum_i \hbar\omega_i (a_i^* a_i + 1/2)$$

Casimir momentum: 2/6

QED of harmonic oscillator in crossed fields



*Conjugate momenta
≠ kinetic momentum*

$$\mathbf{p}_1 = m_1 \mathbf{v}_1 + e \mathbf{A}_0(\mathbf{r}_1)$$
$$\mathbf{p}_2 = m_1 \mathbf{v}_2 - e \mathbf{A}_0(\mathbf{r}_2)$$

*Pseudo momentum is
conserved*

$$\hat{\mathbf{K}} = \mathbf{p}_1 + \mathbf{p}_2 + \frac{1}{2} e \mathbf{B}_0 \times \mathbf{r}_{21} = \mathbf{P}_{kin} + e \mathbf{B}_0 \times \mathbf{r}$$

$$[\mathbf{K}, H] = 0$$

Coulomb Gauge

Ground state changes due to coupling with quantum vacuum

$$\begin{aligned}
 |\tilde{\Psi}_0\rangle = & \left[1 - \frac{1}{2} \sum'_{i\mathbf{Q}_n} \frac{|W_{i\mathbf{Q}_n,0\mathbf{Q}_00}|^2}{(E_{0\mathbf{Q}_00} - E_{i\mathbf{Q}_n})^2} \right] |0\mathbf{Q}_0\{0\}\rangle \\
 & + \sum_{i\mathbf{Q}_n} \frac{W_{i\mathbf{Q}_n,0\mathbf{Q}_00}}{E_{0\mathbf{Q}_00} - E_{i\mathbf{Q}_n}} |i\mathbf{Q}_n\rangle \\
 & + \sum'_{i\mathbf{Q}_n} \sum'_{i'\mathbf{Q}'n'} \frac{W_{i\mathbf{Q}_n,i'\mathbf{Q}'n'} W_{i'\mathbf{Q}'n',0\mathbf{Q}_00}}{(E_{0\mathbf{Q}_00} - E_{i\mathbf{Q}_n})(E_{0\mathbf{Q}_00} - E_{i'\mathbf{Q}'n'})} |i\mathbf{Q}_n\rangle \\
 & - W_{0\mathbf{Q}_00,0\mathbf{Q}_00} \sum'_{i\mathbf{Q}_n} \frac{W_{i\mathbf{Q}_n,0\mathbf{Q}_00}}{(E_{0\mathbf{Q}_00} - E_{i\mathbf{Q}_n})^2} |i\mathbf{Q}_n\rangle
 \end{aligned}$$

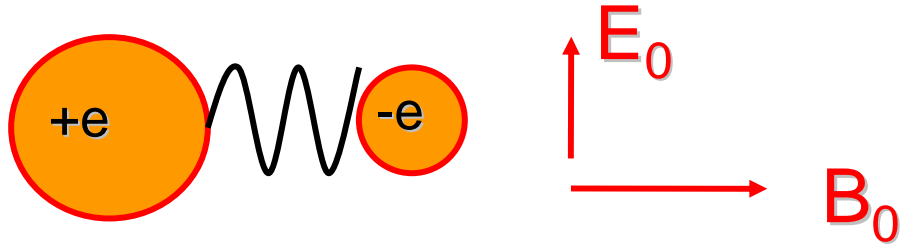
$\delta M_{\mathbf{v}+}$
 $\langle \psi_0 | e\mathbf{A} | \psi_0 \rangle$

$0.84\alpha^3$
 $\alpha(\omega=0, \mu+\delta\mu)$

$\langle 0 | \mathbf{A} | 0 \rangle = 0$

Casimir momentum: 3/6

QED of hydrogen atom in crossed fields



$$\hat{K} = m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2 + e\mathbf{A}(\mathbf{r}_1) - e\mathbf{A}(\mathbf{r}_2) + e\mathbf{B}_0 \times \mathbf{r}_{21} + \sum_i \hbar \mathbf{k}_i (a_i^* a_i + 12)$$

No multipole approximation in $\mathbf{A}(\mathbf{r}) \propto \sum_{gk} \mathbf{g}_k \exp(i\mathbf{k}\mathbf{r}) a_{gk} + c.c$

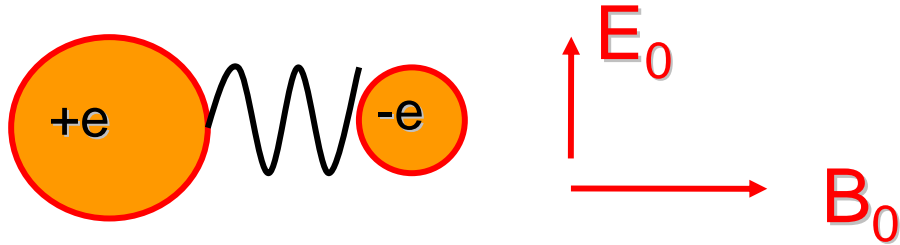
$$\langle \Psi_0 | \mathbf{K} | \Psi_0 \rangle = M\mathbf{v} + \delta M \mathbf{v} + \frac{8}{3} \frac{E_0}{c^2} \mathbf{v}$$

$$+ \varepsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 + \varepsilon_0 \delta\mu \partial_\mu \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 + \mathbf{K}_1 + \mathbf{K}_2$$

$$\delta M = \delta(m_1 + m_2) \quad \delta\mu = \delta\left(\frac{m_1 m_2}{m_1 + m_2}\right) \quad \delta m_i = \frac{4}{3\pi} \alpha \hbar \int_0^\infty dk \frac{\hbar k}{\hbar^2 k^2 / 2m_i + \hbar kc}$$

Casimir momentum: 4/6

QED of hydrogen atom in crossed fields



$$\hat{K} = m_1 \mathbf{v}_1 + m_2 \mathbf{v}_2 + e\mathbf{A}(\mathbf{r}_1) - e\mathbf{A}(\mathbf{r}_2) + e\mathbf{B}_0 \times \mathbf{r}_{21} + \sum_i \hbar \mathbf{k}_i \left(a_i^* a_i + \frac{1}{2} \right)$$

$$-\frac{m_1^2 \mathbf{v}_1^2}{2c_0^2} \mathbf{v}_1 + \text{idem } 2$$

$$\langle \Psi_0 | \mathbf{K} | \Psi_0 \rangle = M\mathbf{v} + \delta M \mathbf{v} + \frac{8 E_0}{3 c_0^2} \mathbf{v} - \frac{5 E_0}{3 c_0^2} \mathbf{v}$$

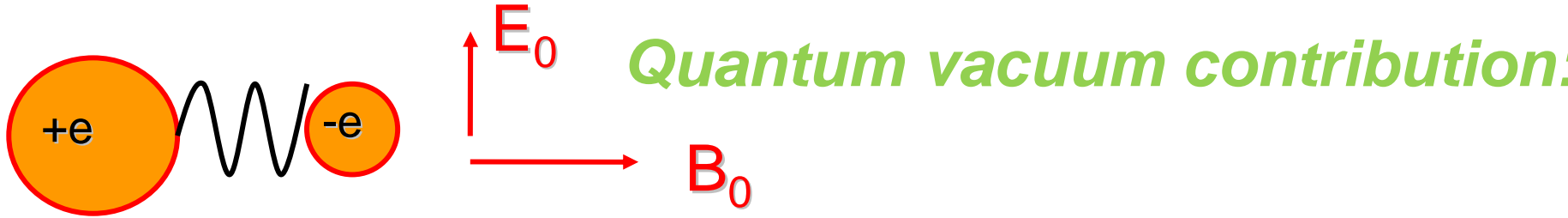
$$+ \varepsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 + \delta(\mu) \varepsilon_0 \partial_\mu \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0$$

$$+ \mathbf{K}_1 + \mathbf{K}_2 + \mathbf{K}_R$$

$$\delta M = \delta(m_1 + m_2) \quad \delta\mu = \delta\left(\frac{m_1 m_2}{m_1 + m_2}\right) \quad \delta m_i = \frac{4}{3\pi} \alpha \hbar \int_0^\infty dk \frac{\hbar k}{\hbar^2 k^2 / 2m_i + \hbar kc}$$

Casimir momentum: 5/6

QED of hydrogen in crossed fields



$$\mathbf{K}_1 = -\mathbf{B}_0 \times \mathbf{E}_0 \frac{1}{3} \frac{e^2 \hbar^2}{a_0 c_0^2 \mu^2} \sum_n \langle 0 | \frac{e^2}{4\pi\epsilon_0 r} \hat{\mathbf{r}} | n \rangle \cdot \frac{1}{(E_n - E_0)^2} \cdot \langle n | \mathbf{r} | 0 \rangle$$
$$= -\epsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 \alpha^2 (0.208 + 0.0045) = -0.21 \alpha^2 \mathbf{K}_A$$

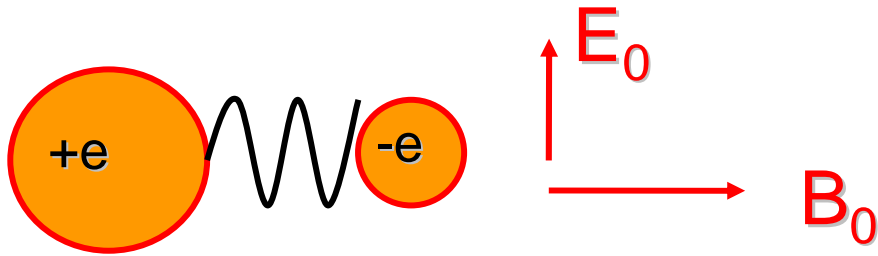
$$\mathbf{K}_2 = +\epsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 \alpha^2 \frac{1}{27} \frac{e^2}{4\pi\epsilon_0 a_0^2} \sum_n \langle 0 | \hat{\mathbf{r}} | n \rangle \cdot \frac{1}{E_n - E_0} \cdot \langle n | \mathbf{r} | 0 \rangle$$
$$= +\epsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0 \alpha^2 (0.079 + 0.018) = +0.1 \alpha^2 \mathbf{K}_A$$

Discrete Rydberg states

Continuous spectrum
assuming
plane waves for electrons

Casimir momentum: 6/6

QED of hydrogen in crossed fields



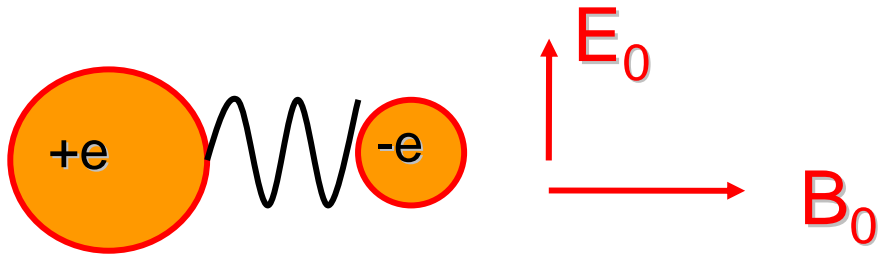
Relativistic contribution:

$$\mathbf{K}_R = -\frac{e^2}{2m_e M c_0^2} \langle 0_E | p^2 (\mathbf{B}_0 \times \mathbf{x}) | 0_E \rangle$$

$$\propto \alpha^2 \frac{m_e}{M} \mathbf{K}_A$$

Casimir momentum: 6/6

QED of hydrogen in crossed fields



$$\langle \mathbf{K} \rangle = \langle \mathbf{K}_{\text{kin}} \rangle + \frac{E_0}{c_0^2} \mathbf{v} + \left(1 - 0.1 \alpha^2\right) \mathbf{K}_A + O(\alpha^3)$$

$$\mathbf{K}_A = \varepsilon_0 \alpha(0) \mathbf{B}_0 \times \mathbf{E}_0$$

**Casimir momentum of H atom exists
and slightly reduces the classical Abraham momentum**

BaVT, Kawka, Rikken, submitted to EPJD

SUMMARY

Casimir momentum in crossed E,B

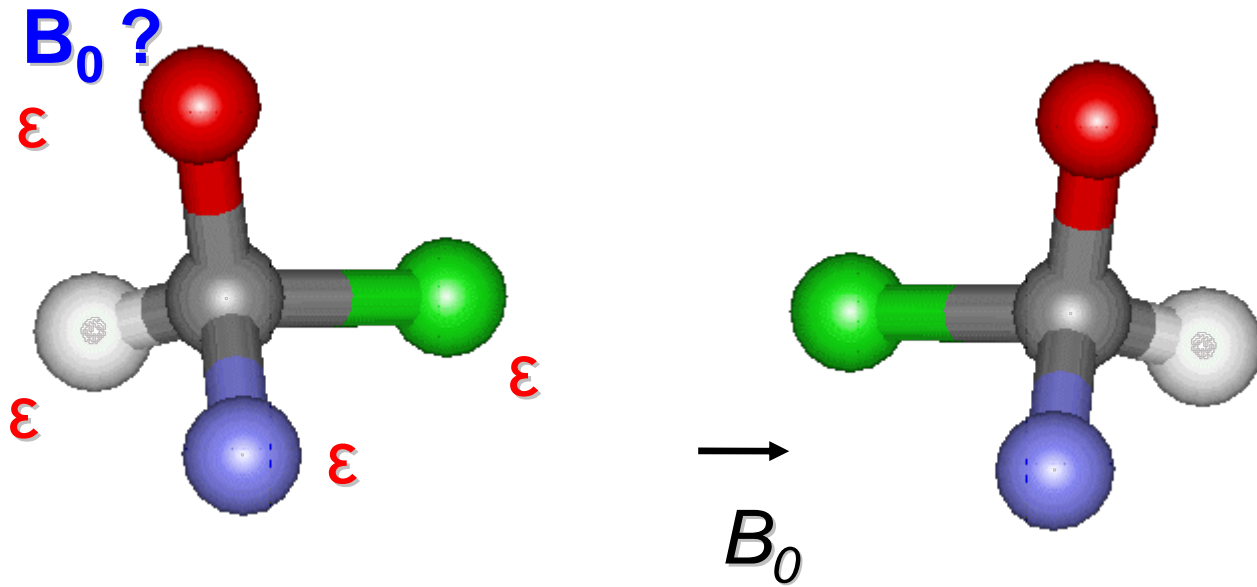
- **Classical Abraham force , linear in E_0 and B_0 , is observed for neutral atoms and for strong dielectrics**
- **QED contribution by Feigel is not observed**
- **UV divergencies disappear in mass renormalization or cancel.**
- **Need to go beyond multipole approximation**
- **Quantum vacuum contributes to Abraham momentum in order $-(1/137)^2$
Will this be $-(Z/137)^2$ for $Z > 1$??**

A Casimir momentum with only magnetic field?

$$\langle \mathbf{E} \times \mathbf{B} \rangle = g\mathbf{B}_0 ?$$

- **Classically no equivalent Abraham version in charge neutral systems**
- **g must be a pseudo scalar**
→ medium must be **chiral (on nanoscale)**
- **Describe chirality microscopically, not phenomenologically**
via «magneto-chiral» index of refraction ($\Delta n = g \mathbf{B}_0 \cdot \mathbf{k}$)
- **Would separate enantiomers using magnetic fields = Pasteurs dream !**
- **Medium must be magnetic since $\langle \mathbf{E} \times \mathbf{H} \rangle = 0$**

Pasteur's dream with a Casimir momentum $\mathbf{P} = \mathbf{g} \times \mathbf{B}_0$



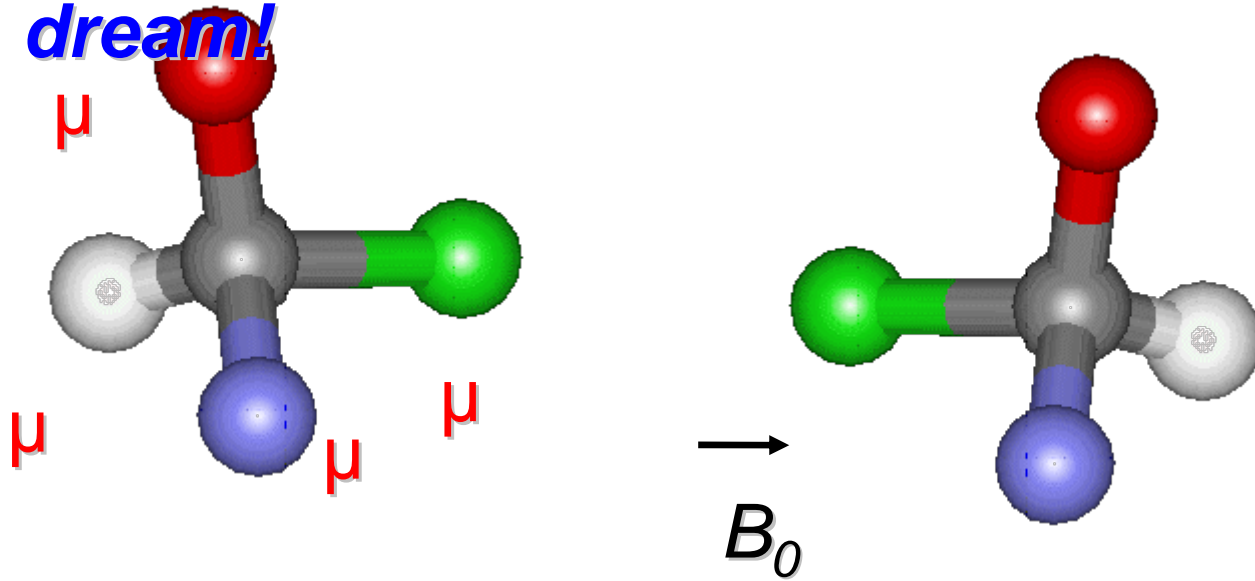
Chiral geometry with electric polarizabilities with Zeeman splitting

$$\alpha(\omega, \sigma) = \frac{4\pi c_0^2}{\omega_0^2} \frac{\gamma}{\omega^2 - \omega_0^2 + i\sigma VB + i\gamma\omega_0}$$



$$\mathbf{B} = \mu_0 \mathbf{H} \Rightarrow \langle 0 | \int d\mathbf{r} \mathbf{E} \times \mathbf{B} | 0 \rangle \propto \langle 0 | \int d\mathbf{r} \mathbf{E} \times \mathbf{H} | 0 \rangle = 0$$

A Casimir momentum $P = g B_0$? Pasteur's dream!



Chiral geometry with magnetic polarizabilities with Zeeman splitting

$$\chi(\omega, \sigma) = \chi(0) \frac{\omega_0^2}{\omega^2 - \omega_0^2 + i\sigma VB + i\gamma\omega}$$

$$\langle 0 | \int d\mathbf{r} \mathbf{E} \times \mathbf{H} | 0 \rangle = 0$$

$$\langle 0 | \int d\mathbf{r} \mathbf{E} \times \mathbf{B} | 0 \rangle = g \mathbf{B}_0$$

$$g = \left(\frac{4\hbar c}{3e\mu_0} \right) (\epsilon \chi(0) \mu_0)^5 \left(\frac{1.4 \times 10^4}{L^{14}} \right)$$

Na Tetraeder $L=10 \text{ nm} \rightarrow g/m = 1 \text{ nm/sec/T}$
 Babington, BaVT, EPL 2011

Lorentz Center

Casimir Physics School and Workshop 2012

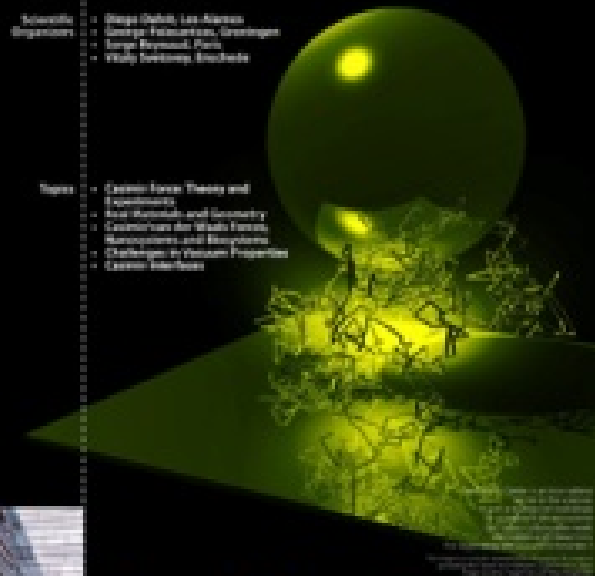
School: 5 - 9 March 2012
Workshop: 12 - 16 March 2012 Leiden, the Netherlands

Scientific Organizers

- Diego Dalvit, Los Alamos
- George Poulakakis, Groningen
- Jorge Reynoso, Paris
- Willy de Nijs, Eindhoven

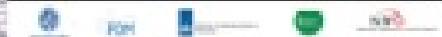
Topics

- Casimir Force Theory and Experiment
- New Materials and Geometry
- Casimir's role in Microfluidics, Nanosystems and Microsystems
- Challenges in Vacuum Properties
- Casimir Interactions



The Casimir effect is a quantum field theory prediction of a force between two uncharged, parallel plates. It is named after the Dutch physicist Hendrik Lorentz, who first proposed it in 1909. The effect is a result of the zero-point energy of the electromagnetic field.

Lorentz Center
 www.lorentzcenter.nl



www.lorentzcenter.nl

Thank you !